

EXACT CONTROLLABILITY OF AN AEROACOUSTIC MODEL

L. COT¹, J.-P. RAYMOND² AND J. VANCOSTENOBLE³

Abstract. We study the exact controllability of a fluid-structure model. The fluctuations of velocity and pressure in the fluid are described by a potential, and the structure is a membrane located in a part Γ_s of the boundary of the domain Ω . The potential ϕ and the transverse displacement z satisfy a coupled system of two wave equations, one in the domain $\Omega \times (0, T)$, the other one in the boundary $\Gamma_s \times (0, T)$. Taking two boundary controls, the first one in a boundary condition satisfied by the potential, and the second one in a boundary condition of the structure equation, we identify the space of controllable initial conditions when the geometrical controllability conditions are satisfied. As in the case of the so-called Helmholtz fluid-structure model [10], the difficulty in the treatment of the observability inequalities, in the definition of very weak solutions, and in the proof of controllability result, comes from the coupling terms of the system. To overcome these difficulties, we show that the variant introduced in [10] of the classical Hilbert Uniqueness Method can be adapted to the aeroacoustic model we consider.

INTRODUCTION

In this paper, we are interested in the control of aeroacoustic models studied in [9], [6], [7], [1], describing acoustic wave propagation in a cavity Ω filled by a fluid flow. In these models, the pressure (the fluctuation from the mean pressure at equilibrium) and the fluctuations of the velocity field are related to a velocity potential ϕ by $p = \rho_f \phi'$ and $v = -\nabla \phi$, where ρ_f is the density of the fluid at equilibrium. The potential ϕ satisfies the undamped wave equation:

$$\phi'' - c_f^2 \Delta \phi = 0 \quad \text{in } \Omega \times (0, T),$$

where c_f^2 is the speed of sound in the fluid. A part Γ_s of the boundary of Ω is occupied by a structure. In [9], [6], [7], [1], and also in the present paper, the structure is a membrane. The structure could also be a beam for a 2D model as in [3], a plate or a shell in a 3D domain [2]. In the following, let us denote by z the transverse displacement of the structure. The common feature of these models is that the continuity condition of velocity on Γ_s is:

$$\partial_n \phi = z' \quad \text{on } \Gamma_s \times (0, T),$$

where $\partial_n \phi$ is the normal derivative of ϕ at the boundary Γ of Ω , and the equation of the transverse displacement of the structure involves

$$-\frac{\rho_f}{\rho_s} \phi'$$

as a source term in the equation. In the case of a membrane the transverse displacement satisfies a wave equation $z'' - c_s^2 \Delta_s z = -\frac{\rho_f}{\rho_s} \phi'$ in $\Gamma_s \times (0, T)$, where c_s^2 is the speed of sound in the structure, ρ_s is the density of the

¹ Laboratoire MIP, UMR CNRS 5640, Université Paul Sabatier, 31062 Toulouse Cedex 4, France, cot@mip.ups-tlse.fr

² Laboratoire MIP, UMR CNRS 5640, Université Paul Sabatier, 31062 Toulouse Cedex 4, France, raymond@mip.ups-tlse.fr

³ Laboratoire MIP, UMR CNRS 5640, Université Paul Sabatier, 31062 Toulouse Cedex 4, France, vancoste@mip.ups-tlse.fr

structure, and Δ_s is the Laplace operator in Γ_s (Γ_s is supposed to be flat). The model is completed by other boundary conditions on $\Gamma \setminus \Gamma_s$ for the potential ϕ , for example $\partial_n \phi = 0$ on a part of $(\Gamma \setminus \Gamma_s) \times (0, T)$ and $\phi = 0$ on an other part of $(\Gamma \setminus \Gamma_s) \times (0, T)$, and a boundary condition on $\partial\Gamma_s$ for the transverse displacement z , for example $z = 0$ on $\partial\Gamma_s$. We would like to control the complete model – including the boundary conditions – by means of a control in the structure equation, or in the boundary conditions of the potential, or by both controls one on the structure and the other one on the fluid. The case of a control on Γ_s in the structure equation has been considered in [9], [7] and [1]. The results in [9] show that, when Ω is a rectangle, only initial conditions satisfying a very specific decay of Fourier modes may be controlled. An approximate controllability result is stated in [7] for particular initial conditions generating waves localized near the structure and called Stoneley waves. In both references, [9] and [7], the controllability is studied by Fourier analysis. In [1], for very particular domains, a controllability result is proved for initial conditions corresponding to the space of finite energy solutions with a control in $(H^1(0, T; L^2(\Gamma_s)))'$, acting only as a distributed term in the structure equation. For other domains, an other controllability result is proved with both controls, one acting as a distributed term in the structure equation and belonging to $L^2(0, T; H^{-\alpha}(\Gamma_s))$ for some $1/4 \leq \alpha \leq 1/3$, and the other one acting in a Neumann boundary condition for the fluid and belonging to $L^2(0, T; L^2(\Gamma \setminus \Gamma_s))$ (see [1, Theorem 1b]). To the best of our knowledge the case of a control acting only on a part of the boundary $\Gamma \setminus \Gamma_s$ for the fluid, and another one acting in the structure equation (either as a distributed control or a boundary control), seems to be open.

In this paper, we would like to answer the following question:

Can one adapt to the fluid-structure models above described the techniques developed in [8] for the wave equation ?

In other words, let us take a control in the structure equation (or in the boundary conditions of the structure equation), and a control in a part of the boundary conditions of the fluid equation (either in a Neumann boundary condition, or in a Dirichlet boundary condition), can we affirm that we are able to control the coupled system ?

At first sight, the answer seems to be obvious: if we control both equations, it should be easy to control the coupled system. We are going to see that the situation is complicated by the presence of the two coupling terms: z' in the boundary conditions of ϕ , and $-\frac{\rho_f}{\rho_s} \phi'$ in the structure equation. Indeed, denoting by $Z = V \times L^2(\Omega) \times H_0^1(\Gamma_s) \times L^2(\Gamma_s)$ the space of initial conditions of finite energy solutions of the coupled system, we explain in section 2 why we cannot define weak solutions of the coupled system for initial data in $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$. This particular feature has already been met with an other simplified fluid-structure model [10]. To overcome this difficulty, following [10], we have to introduce an other space of initial data, denoted by Y throughout the paper, in order to adapt the Hilbert Uniqueness Method to the fluid-structure model that we consider. Let us now precisely describe the fluid-structure model we want to control. We suppose that Ω is a simply connected bounded domain in \mathbb{R}^2 , its boundary Γ is split into two parts: a flat part Γ_s corresponding to the membrane, and the complementary part Γ_0 on which we impose a Dirichlet boundary condition. The uncontrolled model is the following one:

$$\left\{ \begin{array}{ll} \phi'' - c_f^2 \Delta \phi = 0 & \text{in } \Omega \times (0, T), \\ \phi = 0 & \text{on } \Gamma_0 \times (0, T), \\ \partial_n \phi = z' & \text{on } \Gamma_s \times (0, T), \\ \phi(0) = \phi^0 \quad \text{and} \quad \phi'(0) = \phi^1 & \text{in } \Omega, \\ z'' - c_s^2 \Delta_s z = -\frac{\rho_f}{\rho_s} \phi' & \text{in } \Gamma_s \times (0, T), \\ z(0, t) = z(L, t) = 0 & \text{in } (0, T), \\ z(0) = z^0 \quad \text{and} \quad z'(0) = z^1 & \text{in } \Gamma_s, \end{array} \right. \quad (1)$$

where $\Omega =]0, L[\times]0, \ell[$, $\Gamma_s =]0, L[\times \{0\}$ and $\Gamma_0 = \Gamma \setminus \Gamma_s$.

In this paper, we consider the case where Ω is a rectangular domain, in order to underline the difficulties due to

the fluid-structure interactions. The only difficulty to extend our controllability result to other kinds of domains is in the proof of the multiplier lemma (see Lemma 5.1). In the case of a rectangular domain, the solution is regular enough to justify the computations in the proof of Lemma 5.1. For others domains, these computations have to be justified.

The main objective of this work is to extend the well-known boundary controllability result of the wave equation stated in [8, Chapter 1, Theorem 6.1]. Here, we take two boundary controls: u_1 acting in a boundary condition of the fluid equation (i.e. on a part of Γ_0), and u_2 acting in a boundary condition of the structure equation (i.e. on a part of $\partial\Gamma_s = \{(0,0);(L,0)\}$). In order to obtain sharp results, we do not consider u_1 and u_2 acting everywhere on Γ_0 and on $\partial\Gamma_s$. According to the well-known results concerning the wave equation (the so-called optic ray condition, see [4]), it is natural to consider u_1 acting on Γ_c , where Γ_c is composed of the upper side and the right-hand side of $\partial\Omega$, and u_2 acting on the point $x = (L,0)$ (see figure 2). More precisely, we consider

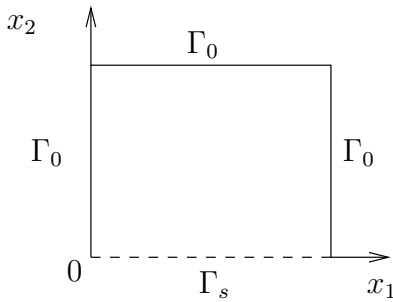


Figure 1

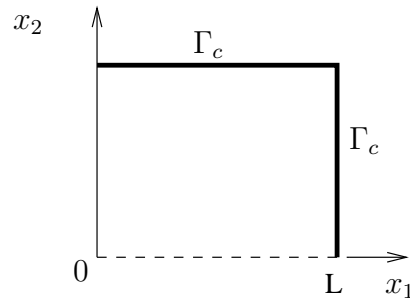


Figure 2

the controlled system:

$$\left\{ \begin{array}{ll} \phi'' - c_f^2 \Delta \phi = 0 & \text{in } \Omega \times (0, T), \\ \phi = u_1 & \text{on } \Gamma_c \times (0, T), \\ \phi = 0 & \text{on } (\Gamma_0 \setminus \Gamma_c) \times (0, T), \\ \partial_n \phi = z' & \text{on } \Gamma_s \times (0, T), \\ \phi(0) = \phi^0 \quad \text{and} \quad \phi'(0) = \phi^1 & \text{in } \Omega, \\ z'' - c_s^2 \Delta_s z = -\frac{\rho_f}{\rho_s} \phi' & \text{in } \Gamma_s \times (0, T), \\ z(0, t) = 0 \quad \text{and} \quad z(L, t) = u_2(t) & \text{in } (0, T), \\ z(0) = z^0 \quad \text{and} \quad z'(0) = z^1 & \text{in } \Gamma_s. \end{array} \right. \quad (2)$$

As already mentioned above the boundary condition $\phi = u_1$, for the fluid equation, may seem to be unrealistic from the physical view point. A more realistic one – a control in a Neumann boundary condition – will be considered in a forthcoming paper, but the analysis will be there more delicate (see [5]). Our main objective is to explain how the Hilbert Uniqueness Method may be adapted to some fluid-structure models. We are going to see that some tools introduced in [10] for a simplified model may be extended to more sophisticated ones. If we summarize, we are interested in the control problem:

Find $T > 0$ and a space of initial conditions E_{ic} such that, for all $(\phi^0, \phi^1, z^0, z^1) \in E_{ic}$, there exists a pair of controls $(u_1, u_2) \in L^2(\Gamma_c \times (0, T)) \times L^2(0, T)$, for which the solution (in a sense to be defined) (ϕ, z) of system (2) obeys:

$$\phi(T) = \phi'(T) = 0 \text{ in } \Omega \quad \text{and} \quad z(T) = z'(T) = 0 \text{ in } \Gamma_s.$$

Accordingly, the paper is organized as follows. Our main results are stated in section 2. We study the well-posedness of the different systems appearing in the Hilbert Uniqueness Method in sections 3 and 4. In section 5, we present the different steps of the Hilbert Uniqueness Method and the proof of the controllability result. In appendix, we establish the three main ingredients, that is, the multiplier lemma, the direct inequality and the inverse inequality.

1. MAIN RESULTS

Throughout the paper, we use the following notation:

$$Q = \Omega \times (0, T), \quad \Sigma = \Gamma \times (0, T), \quad \Sigma_0 = \Gamma_0 \times (0, T), \quad \Sigma_s = \Gamma_s \times (0, T), \quad \Sigma_c = \Gamma_c \times (0, T), \quad \sigma_s = \partial\Gamma_s \times (0, T).$$

1.1. Well-posedness of the homogeneous problem

We introduce the space

$$V = \left\{ \phi \in H^1(\Omega) \mid \phi|_{\Gamma_0} = 0 \right\}.$$

Let us set $\alpha = c_f^2 \frac{\rho_s}{\rho_f}$. The space

$$Z = V \times L^2(\Omega) \times H_0^1(\Gamma_s) \times L^2(\Gamma_s),$$

equipped with the inner product

$$\left((\phi_0, \phi_1, z_0, z_1), (\hat{\phi}_0, \hat{\phi}_1, \hat{z}_0, \hat{z}_1) \right)_Z = c_f^2 \int_{\Omega} \nabla \phi_0 \nabla \hat{\phi}_0 + \int_{\Omega} \phi_1 \hat{\phi}_1 + \alpha c_s^2 \int_{\Gamma_s} \nabla_s z_0 \nabla_s \hat{z}_0 + \alpha \int_{\Gamma_s} z_1 \hat{z}_1,$$

and the associated norm

$$\|(\phi_0, \phi_1, z_0, z_1)\|_Z := \left(c_f^2 \|\phi_0\|_V^2 + \|\phi_1\|_{L^2(\Omega)}^2 + \alpha c_s^2 \|z_0\|_{H_0^1(\Gamma_s)}^2 + \alpha \|z_1\|_{L^2(\Gamma_s)}^2 \right)^{1/2},$$

is a Hilbert space. Observe that $\nabla_s z$ is nothing but $\frac{dz}{dx_1}$. Let us define the energy of regular solutions (ϕ, z) of system (1) by:

$$E(t) := \frac{1}{2} \int_{\Omega} (\phi'^2 + c_f^2 |\nabla \phi|^2) + \frac{\alpha}{2} \int_{\Gamma_s} (z'^2 + c_s^2 |\nabla_s z|^2). \quad (3)$$

Classical computations show that regular solutions (ϕ, z) of (1) satisfy:

$$E(t) = E(0) \quad \text{for all } t \in [0, T].$$

To study system (1) we rewrite it as a first order evolution equation. Setting $U = (\phi, \phi', z, z')$, the system (1) may be rewritten in the form:

$$\frac{dU}{dt} = AU \quad \text{in } (0, T), \quad U(0) = U_0,$$

where

$$AU := A \begin{pmatrix} \phi_0 \\ \phi_1 \\ z_0 \\ z_1 \end{pmatrix} = \begin{pmatrix} \phi_1 \\ c_f^2 \Delta \phi_0 \\ z_1 \\ c_s^2 \Delta_s z_0 - \frac{\rho_f}{\rho_s} \phi_1 \end{pmatrix}, \quad U_0 := \begin{pmatrix} \phi^0 \\ \phi^1 \\ z^0 \\ z^1 \end{pmatrix},$$

and the domain of A in Z is defined by:

$$D(A) = \left\{ (\phi_0, \phi_1, z_0, z_1) \in Z \mid \Delta \phi_0 \in L^2(\Omega), \phi_1 \in V, z_0 \in H^2(\Gamma_s), z_1 \in H_0^1(\Gamma_s), \text{ and } \partial_n \phi_0 = z_1 \text{ on } \Gamma_s \right\}.$$

Let us give a smoothness result, which is crucial in the following to justify all of our computations.

Theorem 1.1. *Let $(\phi_0, \phi_1, z_0, z_1)$ be given in $D(A)$. Then $\phi_0 \in H^2(\Omega)$.*

Proof. Let $(\phi_0, \phi_1, z_0, z_1)$ be given in $D(A)$. Then ϕ_0 satisfies:

$$-\Delta\phi_0 \in L^2(\Omega), \quad \phi_0 = 0 \quad \text{on} \quad \Gamma_0, \quad \partial_n\phi_0 = z_1 \quad \text{on} \quad \Gamma_s.$$

Since z_1 belongs to $H_0^1(\Gamma_s)$ and Ω is a rectangular domain, it can be shown that $\phi_0 \in H^2(\Omega)$. \square

Next, proving that $(A, D(A))$ is the infinitesimal generator of a strongly continuous unitary group (see Lemma 2.1), we can state the following theorem.

Theorem 1.2. *(i) For all $(\phi^0, \phi^1, z^0, z^1) \in Z$, the system (1) admits a unique weak solution $(\phi, \phi', z, z') \in \mathcal{C}([0, T]; Z)$ which satisfies the estimate:*

$$\|(\phi, \phi', z, z')\|_{\mathcal{C}([0, T]; Z)} \leq C \|(\phi^0, \phi^1, z^0, z^1)\|_Z.$$

(ii) If moreover $(\phi^0, \phi^1, z^0, z^1) \in D(A)$, the solution obeys:

$$\|(\phi, \phi', z, z')\|_{\mathcal{C}([0, T]; D(A))} \leq C \|(\phi^0, \phi^1, z^0, z^1)\|_{D(A)}.$$

Remark 1.3. Let us remark that Theorems 1.2 and 1.1 imply that the solution (ϕ, z) of (1) associated with initial data $(\phi^0, \phi^1, z^0, z^1) \in D(A)$ satisfies: $(\phi, z) \in \mathcal{C}([0, T]; H^2(\Omega) \times H^2(\Gamma_s))$.

We would like to extend to system (1) the well known controllability result of the wave equation with a Dirichlet control [8]. For that it should be natural to study the system (1) with initial conditions in $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$. Unfortunately, because of the coupling term $-\frac{\rho_f}{\rho_s}\phi'$ in the equation in $\Gamma_s \times (0, T)$, we cannot define solutions to system (1) for initial conditions in this space.

To extend the notion of solutions to system (1), we can use the extrapolation method. We denote by $(A^*, D(A^*))$ the adjoint of $(A, D(A))$ with respect to the pivot space Z , and by $(D(A^*))^*$ the dual of $D(A^*)$ with respect to the pivot space Z . We verify that $D(A) = D(A^*)$, and we have

$$D(A) \hookrightarrow Z \hookrightarrow (D(A^*))^*$$

with dense and continuous imbeddings. The strongly continuous group $(e^{tA})_{t \in \mathbb{R}}$ on Z may be extended to a strongly continuous group on $(D(A^*))^*$. We still denote it by $(e^{tA})_{t \in \mathbb{R}}$. Following the approach introduced in [10] for the Helmholtz model, we can show that

$$(\phi^0, \phi^1, z^0, z^1) \mapsto \|(\phi^0, \phi^1, z^0, z^1)\|_Y = \left(\|\phi^0\|_{L^2(\Omega)}^2 + \|\phi^1\|_{V'}^2 + \|z^0\|_{L^2(\Gamma_s)}^2 + \|\alpha z^1 + c_f^2 \phi^0\|_{H^{-1}(\Gamma_s)}^2 \right)^{1/2}$$

is a norm on Z equivalent to the norm

$$\|(\phi^0, \phi^1, z^0, z^1)\|_{(D(A^*))^*} = \sup_{(\psi^0, \psi^1, w^0, w^1) \in D(A^*)} \frac{\left((\phi^0, \phi^1, z^0, z^1), (\psi^0, \psi^1, w^0, w^1) \right)_Z}{\|(\psi^0, \psi^1, w^0, w^1)\|_{D(A^*)}}.$$

(Such a result can be proved as in [5]). Thus the closure of Z in this norm $\|\cdot\|_Y$, denoted by Y , can be identified algebraically and topologically with $(D(A^*))^*$. Therefore we can state the following result.

Theorem 1.4. *For all $(\phi^0, \phi^1, z^0, z^1) \in Y$, system (1) admits a unique weak solution (ϕ, z) (in the sense of semigroup theory) which belongs to $\mathcal{C}([0, T]; Y)$ and satisfies*

$$\|(\phi, \phi', z, z')\|_{\mathcal{C}([0, T]; Y)} \leq C \|(\phi^0, \phi^1, z^0, z^1)\|_Y.$$

1.2. Well-posedness of the controllability problem

Since we have to deal with nonhomogeneous boundary conditions in system (2), with controls $(u_1, u_2) \in L^2(\Sigma_c) \times L^2(0, T)$ and with initial conditions in Y , we define weak solutions by the transposition method. To this end we introduce the adjoint system:

$$\begin{cases} \psi'' - c_f^2 \Delta \psi = f & \text{in } Q, \\ \psi = 0 & \text{on } \Sigma_0, \\ \partial_n \psi = w' & \text{on } \Sigma_s, \\ \psi(T) = 0 \quad \text{and} \quad \psi'(T) = 0 & \text{in } \Omega, \\ w'' - c_s^2 \Delta_s w = -\frac{\rho_f}{\rho_s} \psi' + g & \text{in } \Sigma_s, \\ w = 0 & \text{on } \sigma_s, \\ w(T) = 0 \quad \text{and} \quad w'(T) = 0 & \text{in } \Gamma_s, \end{cases} \quad (4)$$

where $(f, g) \in L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$ are given. Notice that, according to the properties of A (see Lemma 2.1) and to Stone's Theorem, the system (4) admits a unique solution (ψ, ψ', w, w') in $\mathcal{C}([0, T]; Z)$.

Now we state two existence and uniqueness results, the first one for initial data in Z , whereas the second one corresponds to initial data in Y .

1.2.1. Weak solutions of (2)

In the case where $(\phi^0, \phi^1, z^0, z^1) \in Z$, we can define weak solutions of system (2) by the transposition method in a classical way.

Definition 1.5. Let $(\phi^0, \phi^1, z^0, z^1)$ belong to Z and (u_1, u_2) belong to $L^2(\Sigma_c) \times L^2(0, T)$. A pair $(\phi, z) \in L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$ is called a weak solution of (2) if and only if, for all $(f, g) \in L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$, the following identity holds:

$$\begin{aligned} \int_Q f \phi + \alpha \int_{\Sigma_s} g z &= \int_\Omega \phi^1 \psi(0) - \int_\Omega \phi^0 \psi'(0) + \int_{\Gamma_s} (\alpha z^1 + c_f^2 \phi^0) w(0) \\ &- \int_{\Gamma_s} z^0 (\alpha w'(0) + c_f^2 \psi(0)) - c_f^2 \int_{\Sigma_c} u_1 \partial_n \psi - \alpha c_s^2 \int_0^T u_2 \partial_{x_1} w(L), \end{aligned} \quad (5)$$

where (ψ, w) is the solution of (4) associated with (f, g) .

Theorem 1.6. Let $(\phi^0, \phi^1, z^0, z^1)$ be given in Z and let (u_1, u_2) be given in $H_0^2(0, T; H^{\frac{3}{2}}(\Gamma_c)) \times H_0^3(0, T)$. The system (2) admits a unique weak solution (ϕ, z) in $L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$ in the sense of definition 1.5. Moreover, $(\phi, \phi', z, z') \in \mathcal{C}([0, T]; H^1(\Omega) \times L^2(\Omega) \times H^1(\Gamma_s) \times L^2(\Gamma_s))$ and it satisfies the estimate

$$\|(\phi, \phi', z, z')\|_{\mathcal{C}([0, T]; Y)} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right\}. \quad (6)$$

1.2.2. Very weak solutions of (2)

Let us observe that if $(\phi^0, \phi^1, z^0, z^1) \in Y$, then neither $\phi^0|_{\Gamma_s}$ nor $\alpha z^1 + c_f^2 \phi^0|_{\Gamma_s}$ are well defined. Thus for initial data in Y , we have to introduce some new tools to define weak solutions of equation (2). We first introduce the bounded linear operator J defined in Z by

$$J(\phi^0, \phi^1, z^0, z^1) = (\phi^0, \phi^1, z^0, \alpha z^1 + c_f^2 \phi^0|_{\Gamma_s}).$$

In Lemma 3.2, we prove that J can be extended to a bounded linear operator \tilde{J} from Y into $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$. Thanks to this new operator \tilde{J} we are now able to define weak solutions of equation with initial data in Y .

Definition 1.7. Let $(\phi^0, \phi^1, z^0, z^1)$ belong to Y , (u_1, u_2) belong to $L^2(\Sigma_c) \times L^2(0, T)$, and set $(\phi^0, \phi^1, z^0, \tilde{z}^1) = \tilde{J}(\phi^0, \phi^1, z^0, z^1)$. A pair $(\phi, z) \in L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$ is called a very weak solution of (2) if and only if, for all $(f, g) \in L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$, the following identity holds:

$$\begin{aligned} \int_Q f\phi + \alpha \int_{\Sigma_s} gz &= \langle \phi^1, \psi(0) \rangle_{V', V} - (\phi^0, \psi'(0))_{L^2(\Omega)} + \langle \tilde{z}^1, w(0) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} \\ &- (z^0, \alpha w'(0) + c_f^2 \psi(0))_{L^2(\Gamma_s)} - c_f^2 \int_{\Sigma_c} u_1 \partial_n \psi - \alpha c_s^2 \int_0^T u_2 \partial_{x_1} w(L), \end{aligned} \quad (7)$$

where (ψ, w) is the solution of (4) associated with (f, g) .

Remark 1.8. Let us observe that if $(\phi^0, \phi^1, z^0, z^1) \in Z$, then $\tilde{z}^1 = \alpha z^1 + c_f^2 \phi^0|_{\Gamma_s}$ and the notion of weak solution in Definition 1.7 coincides with the one of Definition 1.5.

To exhibit some properties of weak solutions of system (2) with initial data in Y , we introduce the following system in which the initial data are not specified:

$$\begin{cases} \psi'' - c_f^2 \Delta \psi = f & \text{in } Q, \\ \psi = 0 & \text{on } \Sigma_0, \\ \partial_n \psi = w' & \text{on } \Sigma_s, \\ w'' - c_s^2 \Delta_s w = -\frac{\rho_f}{\rho_s} \psi' + g & \text{in } \Sigma_s, \\ w = 0 & \text{on } \sigma_s, \end{cases} \quad (8)$$

where $(\psi, \psi', w, w') \in \mathcal{C}([0, T]; Z)$ and $(f, g) \in L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$.

Theorem 1.9. Let $(\phi^0, \phi^1, z^0, z^1)$ be in Y and let (u_1, u_2) be in $L^2(\Sigma_c) \times L^2(0, T)$. The system (2) admits a unique very weak solution (ϕ, z) in $L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$. This solution belongs to $\mathcal{C}([0, T]; L^2(\Omega)) \cap \mathcal{C}^1([0, T]; V') \times \mathcal{C}([0, T]; L^2(\Gamma_s))$ and satisfies the estimate

$$\|\phi\|_{L^\infty(0, T; L^2(\Omega))} + \|\phi'\|_{L^\infty(0, T; V')} + \|z\|_{L^\infty(0, T; L^2(\Gamma_s))} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right\}.$$

Moreover, setting $(\phi^0, \phi^1, z^0, \tilde{z}^1) = \tilde{J}(\phi^0, \phi^1, z^0, z^1)$, there exists a unique function $\tilde{z} \in \mathcal{C}([0, T]; H^{-1}(\Gamma_s))$ satisfying $\tilde{z}(0) = \tilde{z}^1$ and the following formula:

$$\begin{aligned} &\int_0^t \int_\Omega f\phi + \alpha \int_0^t \int_{\Gamma_s} gz \\ &= \langle \phi^1, \psi(0) \rangle_{V', V} - \langle \phi'(t), \psi(t) \rangle_{V', V} - (\phi^0, \psi'(0))_{L^2(\Omega)} + (\phi(t), \psi'(t))_{L^2(\Omega)} \\ &+ \langle \tilde{z}^1, w(0) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} - \langle \tilde{z}(t), w(t) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} - (z^0, \alpha w'(0) + c_f^2 \psi(0))_{L^2(\Gamma_s)} \\ &+ (z(t), \alpha w'(t) + c_f^2 \psi(t))_{L^2(\Gamma_s)} - c_f^2 \int_0^t \int_{\Gamma_c} u_1 \partial_n \psi - \alpha c_s^2 \int_0^t u_2 \partial_{x_1} w(L), \end{aligned} \quad (9)$$

for all t in $[0, T]$ and all $(\psi, \psi', w, w', f, g) \in \mathcal{C}([0, T]; Z) \times L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$ obeying (8). Besides, there exists $C > 0$ independent of $(\phi^0, \phi^1, z^0, z^1)$ such that

$$\|\tilde{z}\|_{L^\infty(0, T; H^{-1}(\Gamma_s))} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right\}. \quad (10)$$

Remark 1.10. If $(\phi, z) \in \mathcal{C}([0, T]; Z)$, then $\tilde{z} = \alpha z' + c_f^2 \phi|_{\Gamma_s}$.

The proof of Theorem 1.9 is based on a generalized direct inequality (see Theorem 5.2 in Appendix).

1.3. Exact controllability

We now state the main result of this paper.

Theorem 1.11. *There exists $T_0 > 0$ such that, for all $T > T_0$, all $(\phi^0, \phi^1, z^0, z^1) \in Y$, there exists $u_1 \in L^2(\Sigma_c)$ and $u_2 \in L^2(0, T)$ for which the solution (ϕ, z) of problem (2) obeys:*

$$\phi(T) = \phi'(T) = 0 \text{ in } \Omega, \quad \text{and} \quad z(T) = z'(T) = 0 \text{ in } \Gamma_s.$$

Remark 1.12. Let us underline that Theorem 1.11 gives a null controllability result. Since $(e^{tA})_{t \in \mathbb{R}}$ is a group on Y , Theorem 1.11 implies the exact controllability in Y .

2. PROOF OF THEOREM 1.2

The proof of Theorem 1.2 is a direct consequence of the following Lemma.

Lemma 2.1. *The operator $(A, D(A))$ is the infinitesimal generator of a strongly continuous unitary group on Z .*

Proof. To prove the lemma, we show that $(A, D(A))$ is m -dissipative and skew-adjoint in Z .

Step 1. $(A, D(A))$ is dissipative in Z . For all $(\phi_0, \phi_1, z_0, z_1)$ in $D(A)$, we have:

$$\begin{aligned} \left(A(\phi_0, \phi_1, z_0, z_1), (\phi_0, \phi_1, z_0, z_1) \right)_Z &= c_f^2 \int_{\Omega} \nabla \phi_1 \nabla \phi_0 + c_f^2 \int_{\Omega} \Delta \phi_0 \phi_1 \\ &+ \alpha c_s^2 \int_{\Gamma_s} \nabla_s z_1 \nabla_s z_0 + \alpha c_s^2 \int_{\Gamma_s} \Delta_s z_0 z_1 - \alpha \frac{\rho_f}{\rho_s} \int_{\Gamma_s} \phi_1 z_1 = 0. \end{aligned}$$

Step 2. $(A, D(A))$ is m -dissipative. Let $F = (f_1, f_2, f_3, f_4)$ be fixed in Z . We want to solve:

$$U = (\phi_0, \phi_1, z_0, z_1) \in D(A), \quad U - AU = F.$$

This is clearly equivalent to

$$\left\{ \begin{array}{ll} \phi_0 \in V, \quad \Delta \phi_0 \in L^2(\Omega), & \phi_0 - c_f^2 \Delta \phi_0 = f_1 + f_2 \text{ in } \Omega, \\ \phi_1 \in V, & \phi_1 = \phi_0 - f_1 \text{ in } \Omega, \\ z_0 \in H^2(\Gamma_s) \cap H_0^1(\Gamma_s), & z_0 - c_s^2 \Delta_s z_0 + \frac{\rho_f}{\rho_s} \phi_0 = \frac{\rho_f}{\rho_s} f_1 + f_3 + f_4 \text{ on } \Gamma_s, \\ z_1 \in H_0^1(\Gamma_s), & z_1 = z_0 - f_3 = \partial_n \phi_0 \text{ on } \Gamma_s. \end{array} \right.$$

If we set

$$g_1 := f_1 + f_2 \in L^2(\Omega), \quad \text{and} \quad g_2 := \frac{\rho_f}{\rho_s} f_1 + f_3 + f_4 \in L^2(\Gamma_s),$$

the system satisfied by (ϕ_0, z_0) may be rewritten in the form

$$\begin{cases} \phi_0 \in V, \Delta\phi_0 \in L^2(\Omega), & \phi_0 - c_f^2 \Delta\phi_0 = g_1 \text{ in } \Omega, \quad \partial_n \phi_0 = z_0 - f_3 \text{ on } \Gamma_s, \\ z_0 \in H^2(\Gamma_s) \cap H_0^1(\Gamma_s), & z_0 - c_s^2 \Delta_s z_0 + \frac{\rho_f}{\rho_s} \phi_0 = g_2 \text{ on } \Gamma_s. \end{cases} \quad (11)$$

Let us equip $W = V \times H_0^1(\Gamma_s)$ with the following inner product:

$$\left((\phi_0, z_0), (\hat{\phi}_0, \hat{z}_0) \right)_W = \int_{\Omega} \phi_0 \hat{\phi}_0 + c_f^2 \int_{\Omega} \nabla \phi_0 \nabla \hat{\phi}_0 + \alpha \int_{\Gamma_s} z_0 \hat{z}_0 + \alpha c_s^2 \int_{\Gamma_s} \nabla_s z_0 \nabla_s \hat{z}_0.$$

We define the bilinear form a on W by:

$$a((\phi_0, z_0), (\hat{\phi}_0, \hat{z}_0)) = \int_{\Omega} \phi_0 \hat{\phi}_0 + c_f^2 \int_{\Omega} \nabla \phi_0 \nabla \hat{\phi}_0 + \alpha \int_{\Gamma_s} z_0 \hat{z}_0 + \alpha c_s^2 \int_{\Gamma_s} \nabla_s z_0 \nabla_s \hat{z}_0 + c_f^2 \int_{\Gamma_s} (\phi_0 \hat{z}_0 - z_0 \hat{\phi}_0),$$

and the linear form L on W by:

$$L(\hat{\phi}_0, \hat{z}_0) = \int_{\Omega} g_1 \hat{\phi}_0 + \alpha \int_{\Gamma_s} g_2 \hat{z}_0 - c_f^2 \int_{\Gamma_s} f_3 \hat{\phi}_0.$$

The existence of a unique weak solution $(\phi_0, z_0) \in W$ to the system (11) follows from the Lax Milgram Lemma applied to the variational problem:

$$\text{find } (\phi_0, z_0) \in W \quad \text{such that} \quad a((\phi_0, z_0), (\hat{\phi}_0, \hat{z}_0)) = L(\hat{\phi}_0, \hat{z}_0), \quad \text{for all } (\hat{\phi}_0, \hat{z}_0) \in W. \quad (12)$$

Next we can easily show that $\Delta\phi_0 \in L^2(\Omega)$ and $z_0 \in H^2(\Gamma_s) \cap H_0^1(\Gamma_s)$. We have $\phi_1 = \phi_0 - f_1 \in V$ and $z_1 = z_0 - f_3 \in H_0^1(\Gamma_s)$. Therefore, there exists $U = (\phi_0, \phi_1, z_0, z_1) \in D(A)$ such that $U - AU = F$, and A is m -dissipative on Z .

Step 3. A is skewadjoint. We can easily check that the adjoint of A with respect to the pivot space Z is defined by

$$D(A^*) = D(A) \quad \text{and} \quad A^* = -A.$$

The proof is complete. \square

3. PROOF OF THEOREMS 1.6 AND 1.9

3.1. Preliminary lemmas

In order to prove Theorems 1.6 and 1.9, we will use some lemmas, which simplify the computations. First, we introduce the system

$$\begin{cases} \psi'' - c_f^2 \Delta\psi = f & \text{in } Q, \\ \psi = 0 & \text{on } \Sigma_0, \\ \partial_n \psi = w' & \text{on } \Sigma_s, \\ \psi(\tau) = \psi_\tau \quad \text{and} \quad \psi'(\tau) = 0 & \text{in } \Omega, \\ w'' - c_s^2 \Delta_s w = -\frac{\rho_f}{\rho_s} \psi' + g & \text{in } \Sigma_s, \\ w = 0 & \text{on } \sigma_s, \\ w(\tau) = w_\tau \quad \text{and} \quad w'(\tau) = 0 & \text{in } \Gamma_s, \end{cases} \quad (13)$$

where $\tau \in [0, T]$, $(\psi_\tau, w_\tau) \in V \times H_0^1(\Gamma_s)$ and $(f, g) \in L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$.

Lemma 3.1. *Let τ be in $[0, T]$, (ψ_τ, w_τ) be in $V \times H_0^1(\Gamma_s)$, and let (f, g) be in $L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$. The system (13) admits a unique weak solution (ψ, ψ', w, w') in $\mathcal{C}([0, T]; Z)$, satisfying the following estimate:*

$$\begin{aligned} & \|(\psi, \psi', w, w')\|_{\mathcal{C}([0, T]; Z)}^2 + c_f^2 \|\partial_n \psi\|_{L^2(\Sigma_c)}^2 + \alpha c_s^2 \|\partial_{x_1} w(L)\|_{L^2(0, T)}^2 \\ & \leq C_T \left\{ \|\psi_\tau\|_V^2 + \|w_\tau\|_{H_0^1(\Gamma_s)}^2 + \|f\|_{L^1(0, T; L^2(\Omega))}^2 + \alpha \|g\|_{L^1(0, T; L^2(\Gamma_s))}^2 \right\}, \end{aligned} \quad (14)$$

where $C_T > 0$ depends on T , but is independent of τ , (ψ_τ, w_τ) and of (f, g) .

Proof. Due to Lemma 2.1 and to classical results for evolution equations, the weak solution of system (13) obeys the estimate

$$\|(\psi, \psi', w, w')\|_{\mathcal{C}([0, T]; Z)}^2 \leq K \left\{ \|\psi_\tau\|_V^2 + \|w_\tau\|_{H_0^1(\Gamma_s)}^2 + \|f\|_{L^1(0, T; L^2(\Omega))}^2 + \alpha \|g\|_{L^1(0, T; L^2(\Gamma_s))}^2 \right\}, \quad (15)$$

where the constant K is independent of τ . From Theorem 5.2 in Appendix we deduce that:

$$c_f^2 \int_{\Sigma_c} |\partial_n \psi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} w(L)|^2 \leq (1 + K) C_T \left\{ \|\psi_\tau\|_V^2 + \|w_\tau\|_{H_0^1(\Gamma_s)}^2 + \|f\|_{L^1(0, T; L^2(\Omega))}^2 + \alpha \|g\|_{L^1(0, T; L^2(\Gamma_s))}^2 \right\}.$$

Since K is independent of τ , taking $C = (1 + K)C_T$, the proof is complete. \square

Recall that J is defined in Z by

$$J(\phi^0, \phi^1, z^0, z^1) = (\phi^0, \phi^1, z^0, \alpha z^1 + c_f^2 \phi^0|_{\Gamma_s}).$$

Lemma 3.2. *J is an isomorphism from Z to Z . It may be uniquely extended to a bounded linear operator \tilde{J} from Y to $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$. Moreover \tilde{J} is an isomorphism from Y to $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$.*

Proof. It is obvious that J is an isomorphism from Z to Z . Observe that J is also continuous from $(Z, \|\cdot\|_Y)$ to $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$, since we have

$$\|J(\phi^0, \phi^1, z^0, z^1)\|_{L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)}^2 = \|(\phi^0, \phi^1, z^0, z^1)\|_Y^2.$$

Thus, J can be uniquely extended to a bounded linear operator \tilde{J} from Y to $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$. Moreover, since

$$\|\tilde{J}(\phi^0, \phi^1, z^0, z^1)\|_{L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)}^2 = \|(\phi^0, \phi^1, z^0, z^1)\|_Y^2,$$

\tilde{J} is an isometry from Y to $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$, and in particular, \tilde{J} is an isomorphism from Y to $L^2(\Omega) \times V' \times L^2(\Gamma_s) \times H^{-1}(\Gamma_s)$. Lemma 3.2 is proved. \square

3.2. Proof of Theorem 1.6

Step 1. *Existence of a solution (ϕ, z) in $L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$. Let $(\phi^0, \phi^1, z^0, z^1)$ be in Z and (u_1, u_2) be in $H_0^2(0, T; H^{\frac{3}{2}}(\Gamma_c)) \times H_0^3(0, T)$. Let \mathcal{L} be the linear form on $L^2(Q) \times L^2(\Sigma_s)$, defined by:*

$$\begin{aligned} \mathcal{L}(f, g) &= \int_\Omega \phi^1 \psi(0) - \int_\Omega \phi^0 \psi'(0) + \int_{\Gamma_s} (\alpha z^1 + c_f^2 \phi^0) w(0) \\ &\quad - \int_{\Gamma_s} z_0 (\alpha w'(0) + c_f^2 \psi(0)) - c_f^2 \int_{\Sigma_c} u_1 \partial_n \psi - \alpha c_s^2 \int_0^T u_2 \partial_{x_1} w(L), \end{aligned}$$

where (ψ, w) is the solution of (4) corresponding to $(f, g) \in L^2(Q) \times L^2(\Sigma_s)$. Let us show that \mathcal{L} is continuous in $L^2(Q) \times L^2(\Sigma_s)$. With Lemma 3.1, we have

$$\begin{aligned} |\mathcal{L}(f, g)| &\leq C \left(\|(\phi_0, \phi_1, z_0, z_1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right) \\ &\quad \times \left(\|(\psi, \psi', w, w')\|_{C([0, T]; Z)} + c_f^2 \|\partial_n \psi\|_{L^2(\Sigma_c)} + \alpha c_s^2 \|\partial_{x_1} w(L)\|_{L^2(0, T)} \right) \\ &\leq C \left\{ \|f\|_{L^2(Q)} + \|g\|_{L^2(\Sigma_s)} \right\}. \end{aligned}$$

Thus \mathcal{L} is continuous in $L^2(Q) \times L^2(\Sigma_s)$ endowed with the inner product

$$\left((f, g), (\phi, z) \right) \longmapsto \int_Q f \phi + \alpha \int_{\Sigma_s} g z.$$

From Riesz' Theorem we deduce that there exists a unique pair (ϕ, z) belonging to $L^2(Q) \times L^2(\Sigma_s)$ such that

$$\mathcal{L}(f, g) = \int_Q f \phi + \alpha \int_{\Sigma_s} g z,$$

for all $(f, g) \in L^2(Q) \times L^2(\Sigma_s)$. Thus the pair (ϕ, z) satisfies identity (5). Now we have to prove that (ϕ, z) belongs to $L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$.

Taking $g = 0$ in formula (5) and using Lemma 3.1, we obtain:

$$\begin{aligned} \left| \int_Q f \phi \right| &\leq \|(\phi^0, \phi^1, z^0, z^1)\|_Y \|(\psi, \psi', w, w')\|_{C([0, T]; Z)} \\ &\quad + c_f^2 \|u_1\|_{L^2(\Sigma_c)} \|\partial_n \psi\|_{L^2(\Sigma_c)} + \alpha c_s^2 \|u_2\|_{L^2(0, T)} \|\partial_{x_1} w(L)\|_{L^2(0, T)} \\ &\leq C \|f\|_{L^1(0, T; L^2(\Omega))} \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right\}. \end{aligned}$$

Therefore, it yields:

$$\|\phi\|_{L^\infty(0, T; L^2(\Omega))} = \sup \left\{ \left| \int_Q f \phi \right|, \|f\|_{L^1(0, T; L^2(\Omega))} = 1 \right\} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right\}.$$

A similar computation shows that:

$$\|z\|_{L^\infty(0, T; L^2(\Gamma_s))} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right\}.$$

Therefore $(\phi, z) \in L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$ is a weak solution of system (2).

Step 2. Regularity of (ϕ, z) . Let $c \in H^3(0, T; H^{\frac{3}{2}}(\Gamma_s))$ be such that

$$\Delta c = 0 \text{ on } \Sigma_s, \quad c(0, t) = 0 \text{ and } c(L, t) = u_2(t) \text{ in } (0, T).$$

(Actually $c(x, t) = x u_2(t)/L$ since Γ_s is a one-dimensional domain). Since $u_1 \in H_0^2(0, T; H^{3/2}(\Omega))$ and $c' \in H^2(0, T; L^2(\Gamma_s))$, there exists $\zeta \in H^2(0, T; H^{3/2-\varepsilon}(\Omega))$ for all $\varepsilon > 0$, such that

$$-\Delta \zeta = 0 \text{ in } Q, \quad \zeta = u_1 \text{ on } \Sigma_c, \quad \zeta = 0 \text{ on } \Sigma_0 \setminus \Sigma_c \quad \text{and} \quad \partial_n \zeta = c' \text{ on } \Sigma_s.$$

(See for example [11].) Let us set $(\Phi, \mathcal{Z}) = (\phi - \zeta, z - c)$. Then, (Φ, \mathcal{Z}) is solution of the following system:

$$\left\{ \begin{array}{ll} \Phi'' - c_f^2 \Delta \Phi = -\zeta'' & \text{in } Q, \\ \Phi = 0 & \text{on } \Sigma_0, \\ \partial_n \Phi = \mathcal{Z}' & \text{on } \Sigma_s, \\ \Phi(0) = \phi^0 - \zeta(0) \quad \text{and} \quad \Phi'(0) = \phi^1 - \zeta'(0) & \text{in } \Omega, \\ \mathcal{Z}'' - c_s^2 \Delta_s \mathcal{Z} = -\frac{\rho_f}{\rho_s} \Phi' - \frac{\rho_f}{\rho_s} \zeta' - c'' & \text{in } \Sigma_s, \\ \mathcal{Z} = 0 & \text{on } \sigma_s, \\ \mathcal{Z}(0) = z^0 - c(0) \quad \text{and} \quad \mathcal{Z}'(0) = z^1 - c'(0) & \text{in } \Gamma_s. \end{array} \right. \quad (16)$$

Hence, setting

$$C = \begin{pmatrix} 0 \\ -\zeta'' \\ 0 \\ -\frac{\rho_f}{\rho_s} \zeta' - c'' \end{pmatrix}, \quad U = \begin{pmatrix} \Phi \\ \Phi' \\ \mathcal{Z} \\ \mathcal{Z}' \end{pmatrix} \quad \text{and} \quad U_0 = \begin{pmatrix} \phi^0 - \zeta(0) \\ \phi^1 - \zeta'(0) \\ z^0 - c(0) \\ z^1 - c'(0) \end{pmatrix},$$

the system (16) can be written in the form:

$$\frac{dU}{dt} = AU + C, \quad U(0) = U_0.$$

From the regularity of (ζ, c) we deduce that $(C, U_0) \in L^2(0, T; Z) \times Z$. Due to Lemma 2.1, the system (16) has a unique solution $(\Phi, \Phi', \mathcal{Z}, \mathcal{Z}') \in \mathcal{C}([0, T]; Z)$, and therefore

$$\phi = \Phi + \zeta \in \mathcal{C}([0, T]; H^1(\Omega)) \cap \mathcal{C}^1([0, T]; L^2(\Omega)), \quad z = \mathcal{Z} + c \in \mathcal{C}([0, T]; H^1(\Gamma_s)) \cap \mathcal{C}^1([0, T]; L^2(\Gamma_s)).$$

Step 3. *Estimate of $\|\phi'\|_{L^\infty(0, T; V')}$.* Fix $\tau \in]0, T[$. Let (ψ, w) be the solution of (13) corresponding to $\psi_\tau \in V$, $w_\tau = 0$ and $(f, g) = (0, 0)$. An easy computation shows that:

$$\begin{aligned} \int_{\Omega} \phi'(\tau) \psi_\tau &= \int_{\Omega} \phi^1 \psi(0) - \int_{\Omega} \phi^0 \psi'(0) + \int_{\Gamma_s} (\alpha z^1 + c_f^2 \phi^0) w(0) \\ &- \int_{\Gamma_s} z^0 (\alpha w'(0) + c_f^2 \psi(0)) + c_f^2 \int_{\Gamma_s} z(\tau) \psi_\tau - c_f^2 \int_0^\tau \int_{\Gamma_c} u_1 \partial_n \psi - \alpha c_s^2 \int_0^\tau \int_{\Gamma_s} u_2 \partial_{x_1} w(L). \end{aligned}$$

Due to Lemma 3.1, there exists $C > 0$, independent of τ , such that:

$$\begin{aligned} \left| \int_{\Omega} \phi'(\tau) \psi_\tau \right| &\leq \|(\phi^0, \phi^1, z^0, z^1)\|_Y \|(\psi, \psi', w, w')\|_{\mathcal{C}([0, T]; Z)} + |c_f^2 \int_{\Gamma_s} z(\tau) \psi_\tau| \\ &\quad + c_f^2 \|u_1\|_{L^2(\Sigma_c)} \|\partial_n \psi\|_{L^2(\Sigma_c)} + \alpha c_s^2 \|u_2\|_{L^2(0, T)} \|\partial_{x_1} w(L)\|_{L^2(0, T)} \\ &\leq C \|\psi_\tau\|_V \left(\|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right) + |c_f^2 \int_{\Gamma_s} z(\tau) \psi_\tau|. \end{aligned}$$

Moreover, due to the estimate of $\|z\|_{L^\infty(0, T; L^2(\Gamma_s))}$ in Step 1, we have

$$\left| \int_{\Gamma_s} z(\tau) \psi_\tau \right| \leq C \|\psi_\tau\|_V \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0, T)} \right\}.$$

Therefore, with the previous estimate we obtain

$$\left| \int_{\Omega} \phi'(\tau) \psi_{\tau} \right| \leq C \|\psi_{\tau}\|_V \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0,T)} \right\},$$

where $C > 0$ is independent of τ . It follows that

$$\|\phi'(\tau)\|_{V'} = \sup \left\{ \left| \int_{\Omega} \phi'(\tau) \psi_{\tau} \right|, \|\psi_{\tau}\|_V = 1 \right\} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0,T)} \right\}.$$

Since the constant C is independent of τ , we obtain:

$$\|\phi'\|_{L^{\infty}(0,T;V')} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0,T)} \right\}.$$

Step 4. Estimate of $\|\alpha z' + c_f^2 \phi\|_{L^{\infty}(0,T;H^{-1}(\Gamma_s))}$. Fix $\tau \in]0, T]$. Let (ψ, w) be the solution of (13) corresponding to $w_{\tau} \in H_0^1(\Gamma_s)$, $\psi_{\tau} = 0$ and $(f, g) = (0, 0)$. An easy computation shows that:

$$\begin{aligned} \int_{\Gamma_s} (\alpha z'(\tau) + c_f^2 \phi(\tau)) w_{\tau} &= \int_{\Omega} \phi^1 \psi(0) - \int_{\Omega} \phi^0 \psi'(0) + \int_{\Gamma_s} (\alpha z^1 + c_f^2 \phi^0) w(0) \\ &- \int_{\Gamma_s} z^0 (\alpha w'(0) + c_f^2 \psi(0)) - c_f^2 \int_0^{\tau} \int_{\Gamma_c} u_1 \partial_n \psi - \alpha c_s^2 \int_0^{\tau} u_2 \partial_{x_1} w(L). \end{aligned}$$

Therefore, we have:

$$\begin{aligned} \left| \int_{\Gamma_s} (\alpha z'(\tau) + c_f^2 \phi(\tau)) w_{\tau} \right| &\leq \|(\phi^0, \phi^1, z^0, z^1)\|_Y \|(\psi, \psi', w, w')\|_{C([0,T];Z)} \\ &+ c_f^2 \|u_1\|_{L^2(\Sigma_c)} \|\partial_n \psi\|_{L^2(\Sigma_c)} + \alpha c_s^2 \|u_2\|_{L^2(0,T)} \|\partial_{x_1} w(L)\|_{L^2(0,T)}. \end{aligned}$$

Due to Lemma 3.1, there exists $C > 0$, independent of τ , such that:

$$\left| \int_{\Gamma_s} (\alpha z'(\tau) + c_f^2 \phi(\tau)) w_{\tau} \right| \leq C \|w_{\tau}\|_{H_0^1(\Gamma_s)} \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0,T)} \right\}.$$

As in step 3, we obtain:

$$\|\alpha z' + c_f^2 \phi\|_{L^{\infty}(0,T;H^{-1}(\Gamma_s))} \leq C \left\{ \|(\phi^0, \phi^1, z^0, z^1)\|_Y + \|u_1\|_{L^2(\Sigma_c)} + \|u_2\|_{L^2(0,T)} \right\}.$$

The proof of Theorem 1.6 is complete. □

3.3. Proof of Theorem 1.9

Let $(\phi^0, \phi^1, z^0, z^1)$ belong to Y , and (u_1, u_2) belong to $L^2(\Sigma_c) \times L^2(0, T)$. Set $(\phi^0, \phi^1, z^0, \tilde{z}^1) = \tilde{J}(\phi^0, \phi^1, z^0, z^1)$. To prove the existence of a weak solution of system (2), we proceed by approximation. There exist sequences $((\phi_n^0, \phi_n^1, z_n^0, z_n^1))_n \subset Z$ and $((u_n^1, u_n^2))_n \subset H^2(0, T; H^{\frac{3}{2}}(\Gamma_c)) \times H_0^3(0, T)$ such that

$$\begin{cases} (\phi_n^0, \phi_n^1, z_n^0) &\rightarrow (\phi^0, \phi^1, z^0) \text{ in } L^2(\Omega) \times V' \times L^2(\Gamma_s), \\ \alpha z_n^1 + c_f^2 \phi_n^0 &\rightarrow \tilde{z}^1 \text{ in } H^{-1}(\Gamma_s), \\ (u_n^1, u_n^2) &\rightarrow (u_1, u_2) \text{ in } L^2(\Sigma_c) \times L^2(0, T). \end{cases} \quad (17)$$

Then, denoting by (ϕ_n, z_n) the solution of (2) associated with $(\phi_n^0, \phi_n^1, z_n^0, z_n^1, u_n^1, u_n^2)$, from Theorem 1.6, it follows that:

$$(\phi_n, \phi'_n, z_n, z'_n) \rightarrow (\phi, \phi', z, z') \text{ in } \mathcal{C}([0, T]; Y), \quad (18)$$

for some $(\phi, \phi', z, z') \in \mathcal{C}([0, T]; Y)$, and (ϕ, ϕ', z, z') satisfies estimate (6). Moreover $(\alpha z'_n + c_f^2 \phi_n)_n$ is a Cauchy sequence in $\mathcal{C}([0, T]; H^{-1}(\Gamma_s))$, and it converges in $\mathcal{C}([0, T]; H^{-1}(\Gamma_s))$ to some $\tilde{z} \in \mathcal{C}([0, T]; H^{-1}(\Gamma_s))$. In particular, we have $\tilde{z}(0) = \tilde{z}^1$. Let (ψ, w) be the solution of (4) corresponding to $(f, g) \in L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$. Let us recall that the pair (ϕ_n, z_n) obeys (see Definition 1.7):

$$\begin{aligned} & \int_Q f \phi_n + \alpha \int_{\Sigma_s} g z_n \\ &= \langle \phi_n^1, \psi(0) \rangle_{V', V} - (\phi_n^0, \psi'(0))_{L^2(\Omega)} + \langle \alpha z_n^1 + c_f^2 \phi_n^0, w(0) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} \\ & - (z_n^0, \alpha w'(0) + c_f^2 \psi(0))_{L^2(\Gamma_s)} - c_f^2 \int_{\Sigma_c} u_n^1 \partial_n \psi - \alpha c_s^2 \int_0^T u_n^2 \partial_{x_1} w(L), \end{aligned} \quad (19)$$

for all $(f, g) \in L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$, where (ψ, w) is the solution of (4). Since $(\partial_n \psi, \partial_{x_1} w(L)) \in L^2(\Sigma_c) \times L^2(0, T)$ and $(\psi, \psi', w, w') \in \mathcal{C}([0, T]; Z)$, we can pass to the limit in the above variational formulation. Thus, the pair (ϕ, z) is a very weak solution of (2). Moreover, we have

$$\begin{aligned} & \int_0^t \int_{\Omega} f \phi_n + \alpha \int_0^t \int_{\Gamma_s} g z_n \\ &= \langle \phi_n^1, \psi(0) \rangle_{V', V} - \langle \phi'_n(t), \psi(t) \rangle_{V', V} - (\phi_n^0, \psi'(0))_{L^2(\Omega)} + (\phi_n(t), \psi'(t))_{L^2(\Omega)} \\ & + \langle \alpha z_n^1 + c_f^2 \phi_n^0, w(0) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} - \langle \alpha z'_n(t) + c_f^2 \phi_n(t), w(t) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} \\ & - (z_n^0, \alpha w'(0) + c_f^2 \psi(0))_{L^2(\Gamma_s)} + (z_n(t), \alpha w'(t) + c_f^2 \psi(t))_{L^2(\Gamma_s)} - c_f^2 \int_0^t \int_{\Gamma_c} u_n^1 \partial_n \psi - \alpha c_s^2 \int_0^t \int_{\Gamma_c} u_n^2 \partial_{x_1} w(L), \end{aligned} \quad (20)$$

for all $t \in [0, T]$ and all $(\psi, \psi', w, w', f, g) \in \mathcal{C}([0, T]; Z) \times L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$ satisfying (8). We obtain (9) by passing to the limit in the above identity.

To prove the uniqueness of the function (ϕ, z, \tilde{z}) , we take $(\phi^0, \phi^1, z^0, z^1) = (0, 0, 0, 0)$, and $(u_1, u_2) = (0, 0)$ in system (2). We first obtain that the corresponding solution (ϕ, z) is equal to zero. Next, identity (9) implies that $\tilde{z} = 0$.

4. PROOF OF THEOREM 1.11

To prove Theorem 1.11, we use the Hilbert Uniqueness Method. If we transpose to system (2) the H. U. Method for the Dirichlet boundary control of the wave equation, we have to consider the solution (θ, d) of the homogeneous problem (1) corresponding to initial conditions $(\theta^0, \theta^1, d^0, d^1)$ in Z .

Next, we have to use $\partial_n \theta|_{\Sigma_c}$ and $\partial_{x_1} d(L)$ as controls in the backward problem:

$$\left\{ \begin{array}{ll} \psi'' - c_f^2 \Delta \psi = 0 & \text{in } Q, \\ \psi = \partial_n \theta & \text{on } \Sigma_c, \\ \psi = 0 & \text{on } \Sigma_0 \setminus \Sigma_c, \\ \partial_n \psi = w' & \text{on } \Sigma_s, \\ \psi(T) = 0 \quad \text{and} \quad \psi'(T) = 0 & \text{in } \Omega, \\ w'' - c_s^2 \Delta_s w = -\frac{\rho_f}{\rho_s} \psi' & \text{in } \Sigma_s, \\ w(L) = \partial_{x_1} d(L) \quad \text{and} \quad w(0) = 0 & \text{on } (0, T), \\ w(T) = 0 \quad \text{and} \quad w'(T) = 0 & \text{in } \Gamma_s. \end{array} \right. \quad (21)$$

In Lemma 4.1, we are going to see that the pairs (θ, d) and (ψ, w) satisfy a formula of the type:

$$\begin{aligned} & -\langle \psi'(0), \theta(0) \rangle_{V', V} + \langle \psi(0), \theta'(0) \rangle_{L^2(\Omega)} - \langle \tilde{w}(0), d(0) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} \\ & + \langle w(0), \alpha d^1 + c_f^2 \theta^0 \rangle_{L^2(\Gamma_s)} = c_f^2 \int_{\Sigma_c} |\partial_n \theta|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} d(L)|^2, \end{aligned} \quad (22)$$

where \tilde{w} is the unique function in $\mathcal{C}([0, T]; H^{-1}(\Gamma_s))$ satisfying (23) (the analogue of (9) for the adjoint system (21)). In the H. U. Method, to control an initial condition $(\phi^0, \phi^1, z^0, z^1) \in Y$, we look for initial data $(\theta^0, \theta^1, d^0, d^1) \in Z$ so that the solution (ψ, w) of (21) obeys:

$$(\psi(0), \psi'(0)) = (\phi^0, \phi^1) \quad \text{in } \Omega, \quad \text{and} \quad (w(0), w'(0)) = (z^0, z^1) \quad \text{on } \Gamma_s.$$

In the classical H. U. Method the existence of such initial data $(\theta^0, \theta^1, d^0, d^1) \in Z$ is obtained by showing that the mapping

$$(\theta^0, \theta^1, d^0, d^1) \longmapsto (\psi'(0), -\psi(0), w'(0), -w(0))$$

is an isomorphism from Z to $Z' \equiv V' \times L^2(\Omega) \times H^{-1}(\Gamma_s) \times L^2(\Gamma_s)$. Due to formula (22), in our case we have to consider the mapping

$$(\theta^0, \theta^1, d^0, d^1) \longmapsto (\psi'(0), -\psi(0), \tilde{w}(0), -w(0)).$$

Henceforth following the method introduced in [10] for the Helmholtz system, we define the operator Λ from Z into Z' by

$$\Lambda(\theta^0, \theta^1, d^0, d^1) = (\psi'(0), -\psi(0), \tilde{w}(0), -w(0)).$$

We are going to show that Λ is an isomorphism from Z to Z' , which allows us to prove our controllability result. For that, in Appendix, we prove a direct and an inverse inequality (see Theorem 5.3 and 5.4), which implies that the mapping

$$(\theta^0, \theta^1, d^0, d^1) \longmapsto \left(c_f^2 \int_{\Sigma_c} |\partial_n \theta|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} d|^2(L) \right)^{1/2}$$

is a norm on Z for T chosen large enough.

From Theorem 1.2, we know that the system (1) admits a unique weak solution (θ, θ', d, d') in $\mathcal{C}([0, T]; Z)$.

The definition of weak solution of system (21) may be adapted from the one given in Definition 1.5 or in Definition 1.7 for the system (2). As in Theorem 1.9, we can prove that system (21) admits a unique solution (ψ, w) in $L^\infty(0, T; L^2(\Omega)) \times L^\infty(0, T; L^2(\Gamma_s))$, and that there exists a unique function $\tilde{w} \in \mathcal{C}([0, T]; H^{-1}(\Gamma_s))$

satisfying the following formula:

$$\begin{aligned}
& \int_t^T \int_{\Omega} f \psi + \alpha \int_t^T \int_{\Gamma_s} g w \\
&= \langle \psi'(t), \zeta(t) \rangle_{V', V} - (\psi(t), \zeta'(t))_{L^2(\Omega)} + \langle \tilde{w}(t), s(t) \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} \\
& - (w(t), \alpha s'(t) + c_f^2 \zeta(t))_{L^2(\Gamma_s)} - c_f^2 \int_t^T \int_{\Gamma_c} \partial_n \theta \partial_n \zeta - \alpha c_s^2 \int_t^T \partial_{x_1} d(L) \partial_{x_1} s(L),
\end{aligned} \tag{23}$$

for all $t \in [0, T]$, and all $(\zeta, \zeta', s, s', f, g)$ belonging to $\mathcal{C}([0, T]; Z) \times L^1(0, T; L^2(\Omega)) \times L^1(0, T; L^2(\Gamma_s))$ and obeying (8).

Therefore, setting $t = 0$ in the above identity, we see that (ψ, w) is the unique very weak solution of (2) associated with $u_1 = \partial_n \theta$, $u_2 = \partial_{x_1} d(L)$ and the initial condition $(\psi^0, \psi^1, w^0, w^1) = (\tilde{J})^{-1}(\psi(0), \psi'(0), w(0), \tilde{w}(0))$.

Before proving that Λ is an isomorphism from Z into Z' , we establish the following lemma.

Lemma 4.1. *For all $(\phi^0, \phi^1, z^0, z^1)$ be in Z and all $(\theta^0, \theta^1, d^0, d^1)$ be in Z , we have*

$$\left\langle \Lambda(\theta^0, \theta^1, d^0, d^1), J(\phi^0, \phi^1, z^0, z^1) \right\rangle_{Z', Z} = c_f^2 \int_{\Sigma_c} \partial_n \phi \partial_n \theta + \alpha c_s^2 \int_0^T \partial_{x_1} z(L) \partial_{x_1} d(L),$$

where (ϕ, z) and (θ, d) are the solutions of (1) corresponding respectively to $(\phi^0, \phi^1, z^0, z^1)$ and $(\theta^0, \theta^1, d^0, d^1)$.

Proof. The identity in the Lemma is nothing but the formula (7), in which the roles of (ϕ, z) and (ψ, w) are permuted:

$$\begin{aligned}
& \langle \psi'(0), \phi^0 \rangle_{V', V} - (\psi(0), \phi^1)_{L^2(\Omega)} + \langle \tilde{w}(0), z^0 \rangle_{H^{-1}(\Gamma_s), H_0^1(\Gamma_s)} \\
& - (w(0), \alpha z^1 + c_f^2 \phi^0)_{L^2(\Gamma_s)} - c_f^2 \int_{\Sigma_c} \partial_n \theta \partial_n \phi - \alpha c_s^2 \int_0^T \partial_{x_1} d(L) \partial_{x_1} z(L) = 0.
\end{aligned}$$

□

Up to the end of this section, we suppose that $T > T_0$, where T_0 is defined in Theorem 5.4. We introduce the seminorm in Z defined by

$$\|(\theta^0, \theta^1, d^0, d^1)\|_Z := \left(c_f^2 \int_{\Sigma_c} |\partial_n \theta|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} d(L)|^2 \right)^{1/2},$$

where (θ, d) is the solution of (1) with the initial condition $(\theta^0, \theta^1, d^0, d^1)$. From the direct and inverse inequalities (see Theorems 5.3 and 5.4 in Appendix), it follows that $\| \cdot \|_Z$ is a norm in Z , equivalent to $\| \cdot \|_Z$. We denote by $((\cdot, \cdot))_Z$ the associated inner product:

$$((\theta^0, \theta^1, d^0, d^1), (\phi^0, \phi^1, z^0, z^1))_Z := c_f^2 \int_{\Sigma_c} \partial_n \theta \partial_n \phi + \alpha c_s^2 \int_0^T \partial_{x_1} d(L) \partial_{x_1} z(L),$$

where (ϕ, z) and (θ, d) are the solutions of (1) respectively associated with the initial conditions $(\phi^0, \phi^1, z^0, z^1)$ and $(\theta^0, \theta^1, d^0, d^1)$.

Lemma 4.2. *The operator Λ is an isomorphism from Z into Z' .*

Proof. Let us first show that Λ is continuous from Z into Z' . Fix $(\theta^0, \theta^1, d^0, d^1) \in Z$. Let us denote (θ, d) the associated solution of (1). Due to Lemma 4.1, for all $(\phi^0, \phi^1, z^0, z^1) \in Z$, we have

$$\left\langle \Lambda(\theta^0, \theta^1, d^0, d^1), (\phi^0, \phi^1, z^0, z^1) \right\rangle_{Z', Z} = ((\theta^0, \theta^1, d^0, d^1), J^{-1}(\phi^0, \phi^1, z^0, z^1))_Z,$$

where (ϕ, z) the solution of (1) associated to $(\phi^0, \phi^1, z^0, z^1)$. Hence, using Lemma 3.2 and the equivalence of the norms $\|\cdot\|_Z$ and $\|\|\cdot\|\|_Z$, it follows that

$$\begin{aligned} |\langle \Lambda(\theta^0, \theta^1, d^0, d^1), (\phi^0, \phi^1, z^0, z^1) \rangle_{Z', Z}| &\leq \|\|\theta^0, \theta^1, d^0, d^1\|\|_Z \|\|J^{-1}(\phi^0, \phi^1, z^0, z^1)\|\|_Z \\ &\leq C \|\|\theta^0, \theta^1, d^0, d^1\|\|_Z \|\|\phi^0, \phi^1, z^0, z^1\|\|_Z, \end{aligned}$$

for all $(\phi^0, \phi^1, z^0, z^1) \in Z$. Therefore, Λ is continuous from Z into Z' .

Let us show that Λ is an isomorphism from Z into Z' . Still with Lemma 4.1, we have:

$$\left\langle J^* \Lambda(\theta^0, \theta^1, d^0, d^1), (\phi^0, \phi^1, z^0, z^1) \right\rangle_{Z', Z} = \left((\theta^0, \theta^1, d^0, d^1), (\phi^0, \phi^1, z^0, z^1) \right)_Z,$$

for all $(\theta^0, \theta^1, d^0, d^1), (\phi^0, \phi^1, z^0, z^1)$ in Z . From Riesz's Theorem we deduce that $J^* \Lambda$ is an isomorphism from Z into Z' . Thus, $\Lambda = (J^*)^{-1}(J^* \Lambda)$ is an isomorphism from Z into Z' . \square

Proof of Theorem 1.11. Let $(\phi^0, \phi^1, z^0, z^1)$ be in Y . Then $\tilde{J}(\phi^0, \phi^1, z^0, z^1) = (\phi^0, \phi^1, z^0, \tilde{z}^1)$ belongs to $L^2(\Omega) \times V' \times L^2(\Omega) \times H^{-1}(\Gamma_s)$, and $(\phi^1, -\phi^0, \tilde{z}^1, -z^0) \in Z'$. Since Λ is an isomorphism from Z to Z' , there exists $(\theta^0, \theta^1, d^0, d^1) \in Z$ such that

$$\Lambda(\theta^0, \theta^1, d^0, d^1) = (\phi^1, -\phi^0, \tilde{z}^1, -z^0). \quad (24)$$

Let us denote (θ, d) the weak solution of the homogeneous problem (1) with the initial condition $(\theta^0, \theta^1, d^0, d^1)$. The solution (ψ, w) of the backward problem (21) is the unique very weak solution of (2) associated with $u_1 = \partial_n \theta$, $u_2 = \partial_{x_1} d(L)$ and the initial condition $(\psi^0, \psi^1, w^0, w^1) = (\tilde{J})^{-1}(\psi(0), \psi'(0), w(0), \tilde{w}(0))$. By definition of Λ , we have

$$\Lambda(\theta^0, \theta^1, d^0, d^1) = (\psi^1, -\psi^0, \tilde{w}^1, -w^0).$$

Thus equality (24) implies that

$$\tilde{J}(\psi^0, \psi^1, w^0, w^1) = \tilde{J}(\phi^0, \phi^1, z^0, z^1).$$

Therefore, setting $u_1 = \partial_n \theta$ and $u_2 = \partial_{x_1} d(L)$, (ψ, w) is solution of our controllability problem. \square

5. APPENDIX: TECHNICAL TOOLS

The direct and inverse inequalities are not obvious and their proof are technical. It is based on a lemma obtained by the multiplier method.

5.1. Multiplier Lemma

Consider the following nonhomogeneous problem:

$$\left\{ \begin{array}{ll} \phi'' - c_f^2 \Delta \phi = f & \text{in } Q, \\ \phi = 0 & \text{on } \Sigma_0, \\ \partial_n \phi = z' & \text{on } \Sigma_s, \\ \phi(0) = \phi^0 \quad \text{and} \quad \phi'(0) = \phi^1 & \text{in } \Omega, \\ z'' - c_s^2 \Delta_s z = -\frac{\rho_f}{\rho_s} \phi' + g & \text{in } \Sigma_s, \\ z = 0 & \text{on } \partial \Gamma_s \times (0, T), \\ z(0) = z^0 \quad \text{and} \quad z'(0) = z^1 & \text{in } \Gamma_s. \end{array} \right. \quad (25)$$

Lemma 5.1. *Let $(\phi^0, \phi^1, z^0, z^1)$ be in $D(A)$, $(f, g) \in W^{1,1}(0, T; L^2(\Omega) \times L^2(\Gamma_s))$, and let (ϕ, z) be the corresponding solution of problem (25). Let $q = (q_1, q_2)$ be in $C^1(\bar{\Omega})^2$ and p be in $C^1(\bar{\Gamma}_s)$. Then the following identities hold:*

$$\begin{aligned} & \frac{c_f^2}{2} \int_{\Sigma_0} q \cdot n |\partial_n \phi|^2 \\ &= \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T + \frac{1}{2} \int_Q \operatorname{div}(q) (\phi'^2 - c_f^2 |\nabla \phi|^2) + c_f^2 \sum_{j=1,2} \int_Q \partial_{x_j} \phi \partial_{x_j} q \cdot \nabla \phi \\ & - \frac{1}{2} \int_{\Sigma_s} q \cdot n (\phi'^2 - c_f^2 |\nabla \phi|^2) - c_f^2 \int_{\Sigma_s} z' q \cdot \nabla \phi - \int_Q f q \cdot \nabla \phi, \end{aligned} \quad (26)$$

and

$$\frac{c_s^2}{2} \int_0^T \left[p |\partial_{x_1} z|^2 \right]_0^L = \left[\int_{\Gamma_s} z' p \cdot \nabla_s z \right]_0^T + \frac{1}{2} \int_{\Sigma_s} \frac{dp}{dx_1} (z'^2 + c_s^2 |\nabla_s z|^2) + \frac{\rho f}{\rho_s} \int_{\Sigma_s} \phi' p \cdot \nabla_s z - \int_{\Sigma_s} g p \cdot \nabla_s z. \quad (27)$$

Proof. We first prove (26). Since $(\phi^0, \phi^1, z^0, z^1) \in D(A)$ and $(f, g) \in W^{1,1}(0, T; L^2(\Omega) \times L^2(\Gamma_s))$, (ϕ, z) belongs to $L^\infty(0, T; H^2(\Omega) \times H^2(\Gamma_s))$. This smoothness result allows us to justify all of the computations in the proof below.

Multiplying by $q \cdot \nabla \phi$ the equation satisfied by ϕ we get:

$$\int_Q (\phi'' - c_f^2 \Delta \phi) q \cdot \nabla \phi = \int_Q f q \cdot \nabla \phi. \quad (28)$$

With an integration by parts and a Green formula, we obtain:

$$\begin{aligned} \int_Q \phi'' q \cdot \nabla \phi &= \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T - \int_Q \phi' q \cdot \nabla \phi' = \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T - \frac{1}{2} \int_Q q \cdot \nabla |\phi'|^2 \\ &= \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T + \frac{1}{2} \int_Q |\phi'|^2 \operatorname{div}(q) - \frac{1}{2} \int_{\Sigma_s} q \cdot n |\phi'|^2. \end{aligned} \quad (29)$$

On the other hand, since $\phi \in H^2(\Omega)$, with a Green formula, we can write:

$$\begin{aligned} \int_Q \Delta \phi q \cdot \nabla \phi &= \int_{\Sigma} \partial_n \phi q \cdot \nabla \phi - \int_Q \nabla \phi \cdot \nabla (q \cdot \nabla \phi) \\ &= \int_{\Sigma} \partial_n \phi q \cdot \nabla \phi - \sum_{j=1,2} \int_Q \partial_{x_j} \phi (q \cdot \partial_{x_j} \nabla \phi + \partial_{x_j} q \cdot \nabla \phi) \\ &= \int_{\Sigma} \partial_n \phi q \cdot \nabla \phi - \frac{1}{2} \sum_{j=1,2} \int_Q q \cdot \nabla |\partial_{x_j} \phi|^2 - \sum_{j=1,2} \int_Q \partial_{x_j} \phi \partial_{x_j} q \cdot \nabla \phi \\ &= \int_{\Sigma} \partial_n \phi q \cdot \nabla \phi + \frac{1}{2} \int_Q \operatorname{div}(q) |\nabla \phi|^2 - \frac{1}{2} \int_{\Sigma} q \cdot n |\nabla \phi|^2 - \sum_{j=1,2} \int_Q \partial_{x_j} \phi \partial_{x_j} q \cdot \nabla \phi. \end{aligned} \quad (30)$$

Thus, from (28), (29) and (30), we deduce:

$$\begin{aligned} & \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T + \frac{1}{2} \int_Q \operatorname{div}(q) |\phi'|^2 - \frac{1}{2} \int_{\Sigma_s} q \cdot n |\phi'|^2 - c_f^2 \int_{\Sigma} \partial_n \phi q \cdot \nabla \phi \\ & - \frac{c_f^2}{2} \int_Q \operatorname{div}(q) |\nabla \phi|^2 + \frac{c_f^2}{2} \int_{\Sigma} q \cdot n |\nabla \phi|^2 + c_f^2 \sum_{j=1,2} \int_Q \partial_{x_j} \phi \partial_{x_j} q \cdot \nabla \phi = \int_Q f q \cdot \nabla \phi. \end{aligned} \quad (31)$$

Moreover, since $\phi = 0$ on Γ_0 , we have:

$$\nabla \phi = \partial_n \phi n \quad \text{on } \Gamma_0,$$

and therefore

$$\begin{aligned} & \frac{c_f^2}{2} \int_{\Sigma} q \cdot n |\nabla \phi|^2 - c_f^2 \int_{\Sigma} \partial_n \phi q \cdot \nabla \phi \\ & = \frac{c_f^2}{2} \int_{\Sigma_0} q \cdot n |\partial_n \phi|^2 + \frac{c_f^2}{2} \int_{\Sigma_s} q \cdot n |\nabla \phi|^2 - c_f^2 \int_{\Sigma_0} q \cdot n |\partial_n \phi|^2 - c_f^2 \int_{\Sigma_s} z' q \cdot \nabla \phi \\ & = -\frac{c_f^2}{2} \int_{\Sigma_0} q \cdot n |\partial_n \phi|^2 + c_f^2 \int_{\Sigma_s} \left(\frac{1}{2} q \cdot n |\nabla \phi|^2 - z' q \cdot \nabla \phi \right). \end{aligned} \quad (32)$$

Finally, (26) may be deduced from (31) and (32).

Notice that identity (27) may be obtained in a similar way. The proof of Lemma 5.1 is complete. \square

In the following we denote by C various positive constants independent of T , and by C_T constants depending on T .

5.2. Generalized Direct inequality

In the theorem below we extend to system (25), the hidden regularity result well known for the wave equation.

Theorem 5.2. *There exists $C_T > 0$, only depending of T , such that, for all $(\phi^0, \phi^1, z^0, z^1) \in Z$ and all $(f, g) \in L^1(0, T; L^2(\Omega) \times L^2(\Gamma_s))$, the weak solution (ϕ, z) of (25) satisfies:*

$$c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 \leq C_T \left\{ \|(\phi, \phi', z, z')\|_{L^\infty(0, T; Z)}^2 + \|f\|_{L^1(0, T; L^2(\Omega))}^2 + \alpha \|g\|_{L^1(0, T; L^2(\Gamma_s))}^2 \right\}. \quad (33)$$

Besides, the function $T \mapsto C_T$ is nondecreasing.

Proof. Observe that, for weak solutions $(\phi, \phi', z, z') \in \mathcal{C}([0, T]; Z)$ of (25), the terms $\partial_n \phi|_{\Gamma_c}$ and $\partial_{x_1} z(L)$ are a priori not defined. Therefore, as for the wave equation, we first prove the estimate (33) for classical solutions of (25), that is in the case when $(\phi^0, \phi^1, z^0, z^1) \in D(A)$ and $(0, f, 0, g) \in W^{1,1}(0, T; Z)$. Next $(\partial_n \phi|_{\Gamma_c}, \partial_{x_1} z(L))$ is defined as the limit in $L^2(\Sigma_c) \times L^2(0, T)$ of a sequence $((\partial_n \phi_k|_{\Gamma_c}, \partial_{x_1} z_k(L)))_k$, where (ϕ_k, z_k) is the classical solution of (25) associated with $(\phi_k^0, \phi_k^1, z_k^0, z_k^1) \in D(A)$ and $(0, f_k, 0, g_k) \in W^{1,1}(0, T; Z)$, and where $((\phi_k^0, \phi_k^1, z_k^0, z_k^1))_k$ converges to $(\phi^0, \phi^1, z^0, z^1)$ in Z , $(f_k)_k$ converges to f in $L^1(0, T; L^2(\Omega))$, $(g_k)_k$ converges to g in $L^1(0, T; L^2(\Gamma_s))$.

Let (ϕ, z) be a classical solution of (25). We set $m(x) = x$ for all $x \in \bar{\Omega}$. Notice that there exists $0 < \delta \leq L$ such that

$$m \cdot n \geq \delta \quad \text{on } \Gamma_c, \quad m \cdot n = 0 \quad \text{on } \Gamma_0 \setminus \Gamma_c, \quad m \cdot n = 0 \quad \text{on } \Gamma_s. \quad (34)$$

Let us apply Lemma 5.1 with

$$q = m \quad \text{and} \quad p(x_1) = x_1 \quad \forall x_1 \in [0, L]. \quad (35)$$

Let us successively estimate the terms in the identity (26). Using the fact that $m \cdot n = 0$ on Γ_s , we obtain:

$$\begin{aligned} \left[\int_{\Omega} \phi' m \cdot \nabla \phi \right]_0^T &\leq 2C \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2, \\ \frac{1}{2} \int_Q \operatorname{div}(m) (\phi'^2 - c_f^2 |\nabla \phi|^2) &\leq CT \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2, \\ c_f^2 \sum_{j=1,2} \int_Q \partial_{x_j} \phi \partial_{x_j} m \cdot \nabla \phi &\leq CT \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2, \\ \frac{1}{2} \int_{\Sigma_s} m \cdot n (\phi'^2 - c_f^2 |\nabla \phi|^2) &= 0, \end{aligned}$$

and

$$\int_Q f m \cdot \nabla \phi \leq C \left\{ \|f\|_{L^1(0,T;L^2(\Omega))}^2 + \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2 \right\}.$$

Since $m \cdot n \geq \delta$ on Γ_c and $m \cdot n = 0$ on $\Gamma_0 \setminus \Gamma_c$, from (26) and the above estimates, it follows that:

$$\delta \frac{c_f^2}{2} \int_{\Sigma_c} |\partial_n \phi|^2 \leq C_T \left\{ \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2 + \|f\|_{L^1(0,T;L^2(\Omega))}^2 \right\} - c_f^2 \int_{\Sigma_s} z' m \cdot \nabla \phi. \quad (36)$$

Similarly, we estimate the terms in the identity (27), and we obtain:

$$\begin{aligned} \alpha \left[\int_{\Gamma_s} z' p \cdot \nabla_s z \right]_0^T &\leq 2C \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2, \\ \frac{\alpha}{2} \int_{\Sigma_s} \partial_{x_1} p (z'^2 + c_s^2 |\nabla_s z|^2) &\leq CT \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2, \\ \alpha \left| \int_{\Sigma_s} g p \cdot \nabla_s z \right| &\leq C \left\{ \alpha \|g\|_{L^1(0,T;L^2(\Gamma_s))}^2 + \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2 \right\}. \end{aligned}$$

Since $p(L) = L \geq \delta$ and $p(0) = 0$, it follows that:

$$\delta \frac{\alpha c_s^2}{2} \int_0^T |\partial_{x_1} z(L)|^2 \leq C_T \left\{ \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2 + \alpha \|g\|_{L^1(0,T;L^2(\Gamma_s))}^2 \right\} + c_f^2 \int_{\Sigma_s} \phi' p \cdot \nabla_s z. \quad (37)$$

The estimates (36) and (37) imply that:

$$\begin{aligned} &\delta \left[\frac{c_f^2}{2} \int_{\Sigma_c} |\partial_n \phi|^2 + \frac{\alpha c_s^2}{2} \int_0^T |\partial_{x_1} z(L)|^2 \right] \\ &\leq C \left\{ \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2 + \|f\|_{L^1(0,T;L^2(\Omega))}^2 + \alpha \|g\|_{L^1(0,T;L^2(\Gamma_s))}^2 \right\} + c_f^2 \int_{\Sigma_s} (\phi' p \cdot \nabla_s z - z' m \cdot \nabla \phi). \end{aligned}$$

Moreover, recalling that $\nabla_s z = \partial_{x_1} z$, $z = 0$ on $\partial\Gamma_s$ and $p(x_1) = x_1 = m(x_1, 0)$ for all $x_1 \in [0, L]$, we can write:

$$\begin{aligned} & \int_{\Sigma_s} (\phi' p \cdot \nabla_s z - z' m \cdot \nabla \phi) \\ &= - \int_{\Sigma_s} \phi p \partial_{x_1} z' + \left[\int_{\Gamma_s} \phi p \partial_{x_1} z \right]_0^T - \int_{\Sigma_s} z' p \partial_{x_1} \phi \\ &= \int_{\Sigma_s} (\partial_{x_1} \phi p + \phi \partial_{x_1} p) z' + \left[\int_{\Gamma_s} \phi p \partial_{x_1} z \right]_0^T - \int_{\Sigma_s} z' p \partial_{x_1} \phi \\ &= \int_{\Sigma_s} \phi \partial_{x_1} p z' + \left[\int_{\Gamma_s} \phi p \partial_{x_1} z \right]_0^T \leq C(T+1) \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2. \end{aligned}$$

Therefore, there exists a constant $C_T > 0$ such that the function $T \mapsto C_T$ is nondecreasing, and for which we have:

$$\frac{c_f^2}{2} \int_{\Sigma_c} |\partial_n \phi|^2 + \frac{\alpha c_s^2}{2} \int_0^T |\partial_{x_1} z(L)|^2 \leq C_T \{ \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2 + \|f\|_{L^1(0,T;L^2(\Omega))}^2 + \alpha \|g\|_{L^1(0,T;L^2(\Gamma_s))}^2 \},$$

which proves (33) for classical solutions. \square

5.3. Direct inequality

Theorem 5.3. *For every $T > 0$, there exists $C_T > 0$ such that, for all $(\phi^0, \phi^1, z^0, z^1) \in Z$, the weak solution (ϕ, z) of (1) satisfies:*

$$c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 \leq C_T E(0).$$

Moreover the function $T \mapsto C_T$ is nondecreasing.

Proof. If $(f, g) = (0, 0)$, the energy of the pair (ϕ, z) is constant, that is $E(0) = \frac{1}{2} \|(\phi, \phi', z, z')\|_{L^\infty(0,T;Z)}^2$, and Theorem 5.3 is a direct consequence of Theorem 5.2. \square

5.4. Inverse inequality

The proof of the inverse inequality is based on the classical multiplier $m(x) = x - x_0$ where x_0 is generally chosen such that $x_0 \in \Omega$ or $x_0 \notin \Omega$, according to the area on which the control is applied. In our case, the best choice seems to be $x_0 = (0, 0)$. Moreover, choosing q and p as in (35) simplifies several non classical terms due to the coupling.

Let us remark that $\operatorname{div}(q) = 2$, $\frac{dp}{dx_1} = 1$, and $\partial_{x_j} q_k = \delta_{kj}$ for all $j, k = 1, 2$.

Theorem 5.4. *There exists a constant $T_0 > 0$ such that, for all $T > T_0$ and all weak solution (ϕ, z) of (1), we have the estimate:*

$$(T - T_0)E(0) \leq C \left\{ c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 \right\},$$

where C is independent of T .

Before proving Theorem 5.4, we first state some technical lemmas.

Lemma 5.5. *Let $\beta > 0$ be given. The weak solution (ϕ, z) of (1) obeys:*

$$\begin{aligned}
& \frac{1}{2} \int_Q (\phi'^2 + \beta c_f^2 |\nabla \phi|^2) + \frac{\alpha}{2} \int_{\Gamma_s} (z'^2 + c_s^2 |\nabla_s z|^2) \\
&= \frac{c_f^2}{2} \int_{\Sigma_0} q \cdot n |\partial_n \phi|^2 + \frac{\alpha c_s^2}{2} \int_0^T p(L) |\partial_{x_1} z(L)|^2 - \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T - \alpha \left[\int_{\Gamma_s} z' p \cdot \nabla_s z \right]_0^T \\
& - \frac{1}{2} \int_Q (\phi'^2 - \beta c_f^2 |\nabla \phi|^2) + c_f^2 \int_{\Sigma_s} z' q \cdot \nabla \phi - c_f^2 \int_{\Sigma_s} \phi' p \cdot \nabla_s z.
\end{aligned} \tag{38}$$

Proof. Applying Lemma 5.1 with $(f, g) = (0, 0)$, and using the fact that $\operatorname{div}(q) = 2$, $\partial_{x_j} q_k = \delta_{kj}$ and $q \cdot n = 0$ on Γ_s , we have

$$\begin{aligned}
& \frac{c_f^2}{2} \int_{\Sigma_0} q \cdot n |\partial_n \phi|^2 = \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T + \int_Q (\phi'^2 - c_f^2 |\nabla \phi|^2) + c_f^2 \int_Q |\nabla \phi|^2 \\
& - \frac{1}{2} \int_{\Sigma_s} q \cdot n (\phi'^2 - c_f^2 |\nabla \phi|^2) - c_f^2 \int_{\Sigma_s} z' q \cdot \nabla \phi \\
& = \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T + \int_Q \phi'^2 - c_f^2 \int_{\Sigma_s} z' q \cdot \nabla \phi.
\end{aligned} \tag{39}$$

Using $\frac{dp}{dx_1} = 1$ and $p(0) = 0$, we obtain:

$$\frac{c_s^2}{2} \int_0^T p(L) |\partial_{x_1} z(L)|^2 = \left[\int_{\Gamma_s} z' p \cdot \nabla_s z \right]_0^T + \frac{1}{2} \int_{\Sigma_s} (z'^2 + c_s^2 |\nabla_s z|^2) + \frac{\rho_f}{\rho_s} \int_{\Sigma_s} \phi' p \cdot \nabla_s z. \tag{40}$$

Let $\beta > 0$ be given and let us write:

$$\int_Q \phi'^2 = \frac{1}{2} \int_Q (\phi'^2 - \beta c_f^2 |\nabla \phi|^2) + \frac{1}{2} \int_Q (\phi'^2 + \beta c_f^2 |\nabla \phi|^2).$$

Then, by adding (39) and (40) multiplied by α , and using the above identity in the right hand side of 39, Lemma 5.5 is proved. \square

Lemma 5.6. *Let $\beta > 0$ be given. The weak solution (ϕ, z) of (1) obeys:*

$$\begin{aligned}
& \frac{1}{2} \int_Q (\phi'^2 + \beta c_f^2 |\nabla \phi|^2) + \frac{\alpha}{2} \int_{\Sigma_s} (z'^2 + c_s^2 |\nabla_s z|^2) \\
& \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 + E(0) \right) + \frac{\beta - 1}{2} \int_Q \phi'^2 + c_f^2 \frac{\beta - 2}{2} \int_{\Sigma_s} z' \phi.
\end{aligned} \tag{41}$$

Proof. Let $\beta > 0$ be given. Since $q \cdot n = 0$ on $\Gamma_0 \setminus \Gamma_c$, with Lemma 5.5 we deduce:

$$\begin{aligned}
& \frac{1}{2} \int_Q (\phi'^2 + \beta c_f^2 |\nabla \phi|^2) + \frac{\alpha}{2} \int_{\Gamma_s} (z'^2 + c_s^2 |\nabla_s z|^2) \\
& \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 \right) + \left| \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T + \alpha \left[\int_{\Gamma_s} z' p \cdot \nabla_s z \right]_0^T \right| \\
& - \frac{1}{2} \int_Q (\phi'^2 - \beta c_f^2 |\nabla \phi|^2) + c_f^2 \int_{\Sigma_s} (z' q \cdot \nabla \phi - \phi' p \cdot \nabla_s z).
\end{aligned} \tag{42}$$

Let us estimate the different terms appearing in the right hand side of (42). The second term can be estimated as follows:

$$\left| \left[\int_{\Omega} \phi' q \cdot \nabla \phi \right]_0^T + \alpha \left[\int_{\Gamma_s} z' p \cdot \nabla_s z \right]_0^T \right| \leq CE(0). \quad (43)$$

For the last integral, with an integration by parts in space in the first term and in time in the second term, and with $z = 0$ on $\partial\Gamma_s$, we obtain:

$$\begin{aligned} & \int_{\Sigma_s} (z' q \cdot \nabla \phi - \phi' p \cdot \nabla_s z) = \int_{\Sigma_s} (z' x_1 \partial_{x_1} \phi - \phi' x_1 \partial_{x_1} z) \\ & = - \int_{\Sigma_s} (x_1 \partial_{x_1} z' + z') \phi + \int_{\Sigma_s} \phi x_1 \partial_{x_1} z' - \left[\int_{\Gamma_s} \phi x_1 \partial_{x_1} z \right]_0^T \\ & = - \int_{\Sigma_s} z' \phi - \left[\int_{\Gamma_s} \phi x_1 \partial_{x_1} z \right]_0^T \leq - \int_{\Sigma_s} z' \phi + CE(0). \end{aligned}$$

Therefore, we have

$$\begin{aligned} & \frac{1}{2} \int_Q (\phi'^2 + \beta c_f^2 |\nabla \phi|^2) + \frac{\alpha}{2} \int_{\Gamma_s} (z'^2 + c_s^2 |\nabla_s z|^2) \\ & \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z|^2(L) + E(0) \right) - \frac{1}{2} \int_Q \phi'^2 + \frac{\beta}{2} c_f^2 \int_Q |\nabla \phi|^2 - c_f^2 \int_{\Sigma_s} z' \phi. \end{aligned} \quad (44)$$

Besides, multiplying by ϕ the equation satisfied by ϕ , we get

$$0 = \int_Q (\phi'' - c_f^2 \Delta \phi) \phi = \left[\int_{\Omega} \phi' \phi \right]_0^T - \int_Q \phi'^2 - c_f^2 \int_{\Sigma} \partial_n \phi \phi + c_f^2 \int_Q |\nabla \phi|^2.$$

Thus, with the boundary conditions $\partial_n \phi = z'$ on Γ_s and $\phi = 0$ on Γ_0 , we obtain:

$$\int_Q (\phi'^2 - c_f^2 |\nabla \phi|^2) = \left[\int_{\Omega} \phi' \phi \right]_0^T - c_f^2 \int_{\Sigma_s} \partial_n \phi \phi = \left[\int_{\Omega} \phi' \phi \right]_0^T - c_f^2 \int_{\Sigma_s} z' \phi.$$

Hence

$$c_f^2 \int_Q |\nabla \phi|^2 = \int_Q \phi'^2 + c_f^2 \int_{\Sigma_s} z' \phi - \left[\int_{\Omega} \phi' \phi \right]_0^T \leq \int_Q \phi'^2 + c_f^2 \int_{\Sigma_s} z' \phi + CE(0). \quad (45)$$

The lemma follows from (44) and (45). \square

Proof of Theorem 5.4. Taking $\beta = 2$ in (41), we first obtain

$$c_f^2 \int_Q |\nabla \phi|^2 + \frac{\alpha}{2} \int_{\Sigma_s} (z'^2 + c_s^2 |\nabla_s z|^2) \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 + E(0) \right). \quad (46)$$

In the same way, taking $\beta = 1$ in (41) and next using (46), it follows that

$$\begin{aligned} TE(0) & \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 + E(0) \right) - \frac{c_f^2}{2} \int_{\Sigma_s} z' \phi \\ & \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 + E(0) + \int_{\Sigma_s} z'^2 + \int_Q |\nabla \phi|^2 \right) \\ & \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 + E(0) \right). \end{aligned}$$

Thus there exists $T_0 > 0$ such that

$$(T - T_0)E(0) \leq C \left(c_f^2 \int_{\Sigma_c} |\partial_n \phi|^2 + \alpha c_s^2 \int_0^T |\partial_{x_1} z(L)|^2 \right).$$

Therefore Theorem 5.4 is established. \square

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